

Introduction: projective-unitary invariance

- Which properties of a set of quantum states are invariant under unitary transformations?

Projective-unitary (PU)-invariant properties

Example of PU-invariant: two-state overlap $\Delta_{AB} = |\langle A|B \rangle|^2 = \text{Tr}(\rho_A \rho_B)$

- Results by Chien and Waldron [1] for pure states:
 - All PU-invariant properties are functions of k -state Bargmann invariants [2]:

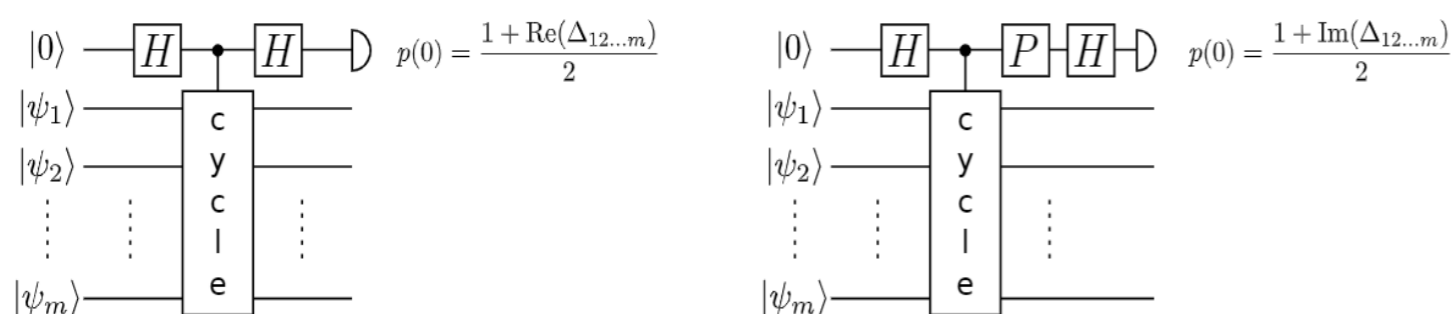
$$\Delta_{ABC\dots K} = \langle A|B \rangle \langle B|C \rangle \langle C|D \rangle \dots \langle K|A \rangle$$

- For N states, we may need up to N -state Bargmann invariants.
- If no pair of states are orthogonal, 3-state invariants suffice for complete PU characterization.
- Bargmann invariants have been discussed in the context of geometric phases, and characterization of multi-photon indistinguishability.

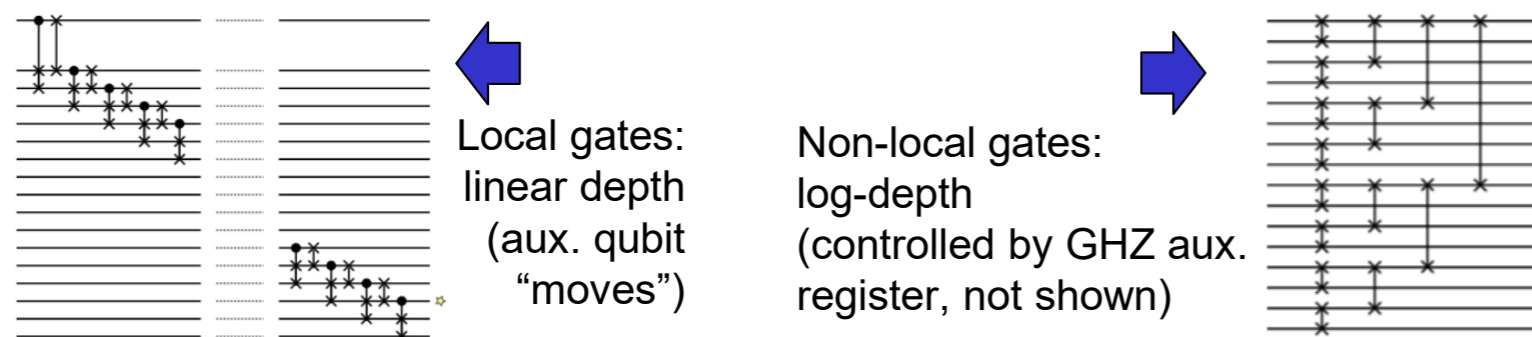
How to measure Bargmann invariants

- The 2-state invariant (overlap) can be measured using the well-known SWAP test:

- We propose the **cycle test**, a generalization to measure real and imaginary parts of Bargmann invariants:



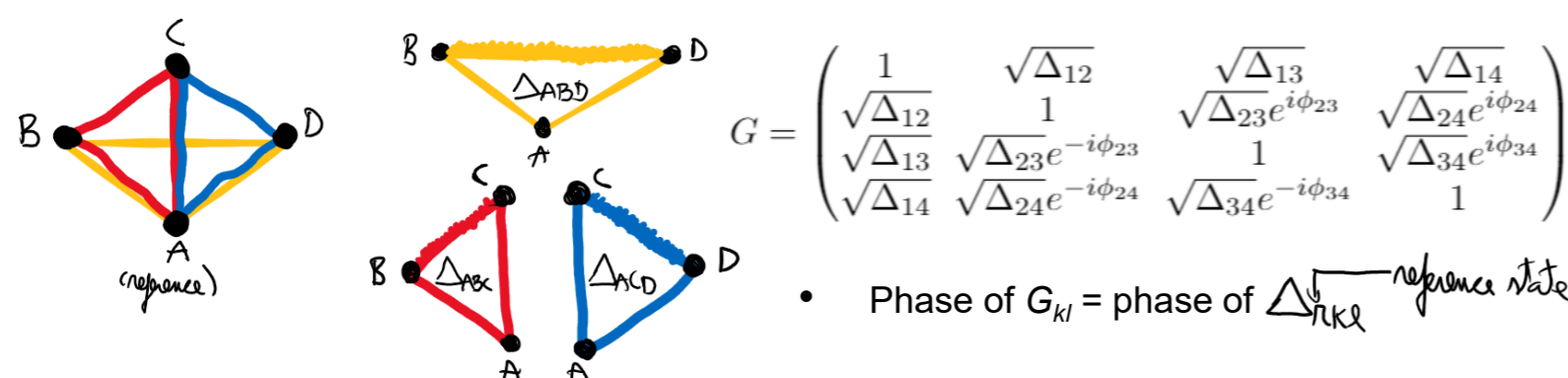
- Circuit on left: used in [3] to evaluate nonlinear functionals of a state.
- Efficient controlled-cycle gate decompositions:



Gram matrix encodes all PU invariants

- PU properties of N generic pure states depend only on 3-state invariants of all pairs of states with a single reference state. Representation: **Gram matrix** G of inner products $G_{ij} \equiv \langle i|j \rangle$

Example: 4 states characterized by 3 3-state invariants, all including reference state A:



- Phase of G_{kl} = phase of Δ_{kkl} (reference state)

- All parameters in G are gauge-invariant and can be measured with cycle tests
- Robustness results: continuity for mixed states with high purity; approximate transformations from one set to another with “close” invariants

Application: Linear independence test

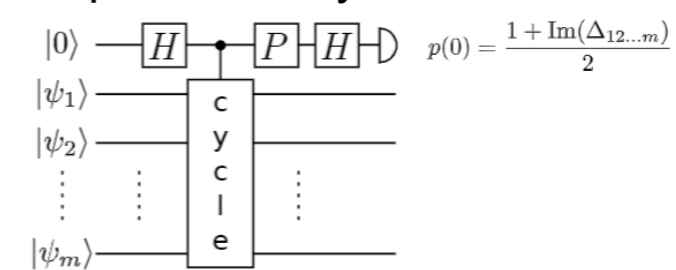
- N states are LI iff $\det(G) > 0$
- Example: $N=3$ states

$$G = \begin{pmatrix} 1 & |\langle \psi_1 | \psi_2 \rangle| & |\langle \psi_1 | \psi_3 \rangle| \\ |\langle \psi_1 | \psi_2 \rangle| & 1 & \langle \psi_2 | \psi_3 \rangle \\ |\langle \psi_1 | \psi_3 \rangle| & \langle \psi_2 | \psi_3 \rangle^* & 1 \end{pmatrix} = \begin{pmatrix} 1 & \sqrt{\Delta_{12}} & \sqrt{\Delta_{13}} \\ \sqrt{\Delta_{12}} & 1 & \sqrt{\Delta_{23}} e^{i \phi_{23}} \\ \sqrt{\Delta_{13}} & \sqrt{\Delta_{23}} e^{-i \phi_{23}} & 1 \end{pmatrix}$$

$$\det(G) > 0 \Leftrightarrow 1 - (\Delta_{12} + \Delta_{13} + \Delta_{23}) + 2\sqrt{\Delta_{12}\Delta_{13}\Delta_{23}} \cos(\phi_{23}) > 0.$$

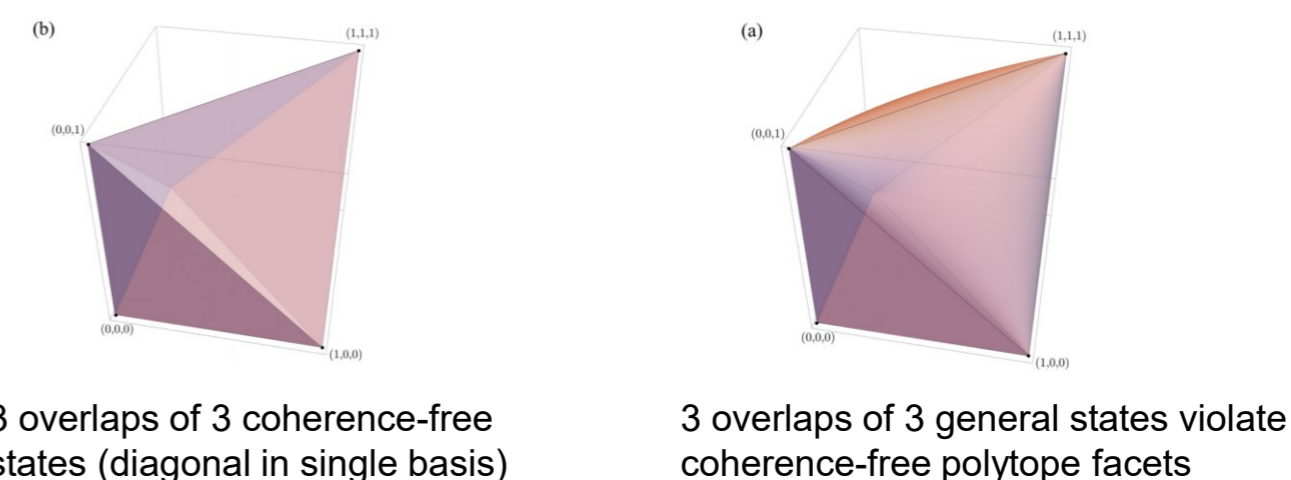
Application: Basis-independent imaginarity witness

- Imaginarity**: necessity of complex amplitudes in quantum theory [4]
- Imaginary part of invariants due to complex amplitudes in input states, measured in a basis-independent way



Application: Basis-independent coherence witnesses

- Bargmann invariant measurements can witness coherence in a basis-independent way
- Example: overlaps of coherence-free (diagonal) states satisfy linear bounds, which are violated by (coherent) states [5]



Conclusion

- We've shown how to measure projective-unitary (Bargmann) invariants of a set of states, encoding a complete set in the Gram matrix G , and applications
- Open questions: algorithmic use in e.g. dimensionality reduction; application in photonic indistinguishability criteria; other foundational relevance (besides witnessing imaginarity and coherence).

References

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